Mn₂ bis(pentalene): a mixed-spin bimetallic with two extremes of bonding within the same molecule[†]

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Structural, magnetic and theoretical studies show that the bimetallic pentalene complex, $Mn_2(C_8H_4^{1,4-Si^3Pr_3})_2$, contains both high and low spin Mn(II) in two very different sites.

Manganese(II) is often shown to be atypical of the first row series, owing to the extra stabilisation energy of the unpaired high-spin d⁵ configuration. In organometallic compounds variable multihaptic bonding of cyclopentadienyl is well known, a feature illustrated most effectively by manganocene $Mn(C_5H_5)_2$.¹ It is unique as the only high-spin metallocene, causing myriad phenomena, as it forms conducting solutions in THF, undergoes facile ligand exchange, and is found in multiple forms in solid and solution states.^{2–5} Molecular manganocene displays a partial spin-crossover at low temperatures, and further work substituting the rings with bulky or electron-donating substituents has created a whole series of both partial and complete spin-crossover compounds.⁶⁻⁹ $Mn(C_5Me_5)_2$ is a low-spin metallocene, with shorter Mn–C bonds, lower reactivity and no observed ligand exchange.10,11 Recent work has also explored the supramolecular chemistry of the [Mn(C₅H₅)₃]⁻ anion.^{12,13}

Metallocene-type environments may also be created using the eight-membered carbocyclic ligand pentalene $[C_8H_6]^{2-}$, which may be commonly considered to act structurally as two fused cyclopentadienyl rings.¹⁴ Katz *et al.* first synthesised the M₂bis(pentalene) dimers for M = Co and Ni, in which twin metallocene-type environments were proposed.^{14–16} Further work has extended this class of compounds to Mo, Rh and Pd centres using a triisopropylsilyl substituted pentalene, in which $\eta^5:\eta^5$ -ligation is found in the crystal structures.^{17–20}

It is clear that the manganocene and M₂bis(pentalene) systems both present striking structural and electronic phenomena, which may lead, for example, to the development of novel magnetic materials. Therefore we have synthesised Mn₂(C₈H₄^{1,4-SiPr₃})₂ and subsequently investigated its magnetic and structural properties. Reaction of an equimolar quantity of C₈H₄^{1,4-SiPr₃}[K]₂ with MnCl₂ in THF gives the bimetallic Mn₂(C₈H₄^{1,4-SiPr₃})₂ (1) as a dark maroon, air-sensitive crystalline solid in 24% yield after workup (eqn (1)).[‡]



A single-crystal X-ray diffraction study reveals an asymmetric coordination environment in the molecule, with two distinctly different manganese centres bound to the ligands in different modes (Fig. 1).§ The structure also lies on a twofold rotation axis through a plane between the opposing pentalene ligands and bisecting the metal centres. Mn-C distances are given in Table 1. One Mn^{II} centre may be considered to be bound in an n⁵:n⁵fashion, with a slighter shorter Mn-C(wingtip) distance (Mn1-C2) than the corresponding bridgehead distances (Mn1-C4, Mn1-C5). The other metal centre is however found significantly closer to the carbon with silvl substituents (Mn2-C8) than the remaining carbon atoms C4-C7 (the next closest is Mn2-C7). Consequent examination of C-C distances in the pentalene skeleton reveals a shorter bond between ring carbons C6-C7 of 1.383(7) Å, compared with an average of 1.444(7) Å for other C-C ring distances, consistent with a localised double bond.¹⁹ The two ligands are slightly tilted away from each other (11°) and significantly twisted from an eclipsed conformation by 24° compared with 10° in Mo₂(η^5 : η^5 -C₈H₄^{1,4-SiⁱPr₃})₂.¹⁷ A disproportionately large displacement ellipsoid for the Mn2 centre is observed, suggestive of a static or dynamic disorder between spatial positions of the metal cation. The Mn(1)...Mn(2) distance of 2.609(2) Å



Fig. 1 Crystal structure of 1 (H and ⁱPr groups removed for clarity).

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	Mn1–C1	Mn1–C2	Mn1–C3	Mn1–C4	Mn1–C5	Mn2–C4	Mn2–C5	Mn2–C6	Mn2–C7	Mn2–C8
Exptl. Calc. $(R = H)$ Calc. $(R = Si^{i}Pr_{3})$ OP $(R = H)$	2.096(5) 2.10 2.13 0.105	2.085(5) 2.07 2.15 0.054	2.225(5) 2.13 2.18 0.124	2.194(5) 2.15 2.27 0.066	2.225(5) 2.18 2.25 0.058	2.471 2.62 2.52 0.030	2.804 2.78 2.85 -0.013	2.692 2.57 2.79 0.031	2.313(5) 2.28 2.42 0.118	2.154(5) 2.25 2.14 0.155

Table 1 Experimental (Exptl.) and calculated (Calc.) Mn-C distances (Å) and atom-atom overlap populations (OP) for Mn₂(C₈H₄R₂₎₂

may be considered short enough to constitute a formal metalmetal bond (which has been recently described between Mn^{II} centres at an internuclear distance of *ca.* 2.7–2.8 Å²¹) and this possibility is discussed following analysis of magnetic data and within a DFT study.

The complex is found to be highly paramagnetic, to the extent that it could not be characterised by NMR methods. Variable-temperature magnetic measurements in solution (Evans method, d_8 -toluene) show that 1 does not obey the Curie–Weiss law; values of effective moment per dimer vary from 7.75 μ_B at 183 K, and decrease to 6.40 μ_B as the temperature was raised to 363 K (Fig. 2(a)). Data obtained for the solid-state susceptibilities (SQUID) meanwhile gave a lower moment of 5.15 μ_B at 4 K increasing slightly to 5.32 μ_B at 300 K (Fig. 2(b)). Fitting these values to a simple Curie–Weiss relationship gave a Curie constant of C = 3.40 emu K mol⁻¹ and a Weiss constant of $\theta = -0.090$ K, showing essentially no additional coupling within the ground state. These values are consistent with an equilibrium between states with



Fig. 2 Plot of molar susceptibility and effective magnetic moment for 1 measured using (a) the Evans' method and (b) a SQUID.

S = 2 and 3, the ground state differing between the solid state (S = 2; $\mu_{spin-only} = 4.90 \ \mu_{B}$) and solution (S = 3; $\mu_{spin-only} = 6.92 \ \mu_{B}$). In the latter case the ground state (S = 3) appears to be thermally depopulated into a state of lower spin multiplicity (S = 2) at higher temperatures.

In order to gain further insight into the bonding and spin-states of this unusual bimetallic, a DFT study was carried out on the model complex $[Mn_2(C_8H_6)_2]$ (2) using the hybrid B3LYP functional.¶ Comparing the relative energies of different spinstates of 2 shows that the septet (S = 3) spin-state is the lowest in energy but the quintet spin-state (S = 2) lies only 4.7 kJ mol⁻¹ higher. The low energy difference between the septet (S = 3) and the quintet (S = 2) spin-states explain the variation of the magnetic moment of 1 in solution upon variation of the temperature. The triplet (S = 1) and the singlet (S = 0) spin-state are 135.3 and 244.9 kJ mol⁻¹, respectively, higher in energy whereas the nonet (S = 4) spin-state is 27.3 kJ mol⁻¹ higher in energy than the septet spin-state.

Geometry optimization of **2** in the septet spin-state also reproduces best the geometric parameters of **1** (Table 1). However, the optimization in a quintet state gave a very similar structure. The asymmetric coordination of the pentalene units to the two different manganese atoms is retained. However, the ring 'slipping' in the optimized geometry of **2** is less accentuated than in the experimental structure of **1**. A QM/MM calculation on **1** with the SiⁱPr₃ groups in the MM region gave a better agreement with experiment (Table 1), thus we attribute part of the twist to steric factors. The preferred binding of Mn2 to C8 is not only evident in the short distance but also the high value between these two atoms in the overlap population matrix (Table 1). This also indicates significant bonding between Mn2 and C7. The long distance between Mn2 and C6 is attributed to the single occupancy of the highest α spin orbital (Fig. 3) which shows a strong antibonding



Fig. 3 The highest energy α SOMO viewed along the Mn–Mn C_2 axis.

interaction between this pair of atoms. The calculated Mn–Mn distance in **2** is slightly longer (2.69 Å) than in the experimental structure of **1** (2.609(2) Å).

An analysis of the atomic spin densities in the septet spin-state (S = 3) for both manganese atoms shows, that *ca.* 1.5 unpaired electrons are localised on Mn1 whereas on Mn2 there are 4.5 unpaired electrons. Almost the same atomic spin density distribution was found in the quintet (S = 2) spin-state, however in this spin-state antiferromagnetic coupling between the two Mn is observed.

To quantify the interaction between the two manganese atoms the Mn–Mn fractional bond order was calculated, by analyzing the interaction of a Mn₂ dimer with the two pentalene units.|| Considering the Mn₂ fragment σ_g , π_u and δ_g orbitals as bonding and σ_u , π_g and δ_u as antibonding and summing up the orbital populations in the complex a fractional bond order of 0.60 can be calculated. This suggests only a weak interaction between the two manganese centres.

Overall DFT suggests that **2** can best be described as having a S = 3 ground state containing a high spin as well as a low spin manganese.

In conclusion, $Mn_2(C_8H_4^{1,4-Si'Pr_3})_2$ has two weakly interacting Mn(II) centres, one high spin and one low-spin, in different coordination environments. The low spin centre shows typical metallocene coordination whereas the high spin centre has a lower coordination number. This contrast strongly with the Mo analogue which is symmetric and shows strong Mo–Mo bonding.

Notes and references

 \ddagger All procedures were carried out under Ar using standard Schlenk techniques in dry solvents. A THF solution of $C_8H_4^{-1.4.5RPr_3}[K]_2$ (246 mg, 0.5 mmol) was added dropwise to a suspension of MnCl₂ (63 mg, 0.5 mmol) in THF. A dark red solution resulted after 24 h of stirring, whereupon all volatiles were removed in vacuo. The dark residue was extracted with toluene (2 \times 40 ml) and filtered through Celite to give a dark maroon red solution. The solution was concentrated (30 ml) and cooling to -50 °C overnight gave the product as dark red crystals. Yield: 110 mg (0.12 mmol, 24% yield). Anal. Calc. for C52H92Si4Mn2: C, 66.48; H, 9.87. Found: C, 66.52; H, 9.76%. MS (EI): m/z = 938 (100%, M⁺); 884 (10%, M⁺ - Mn). § Crystal data for 1: M = 939.50, monoclinic, space group C2/c (no. 15), a =20.5816(9), b = 20.9291(11), c = 12.3807(4) Å, β = 94.419(3)°, V = 5317.2(4) Å³, D_c = 1.17 Mg m⁻³, T = 173(2) K, λ(Mo-Kα) = 0.71073 Å, Z = 4, μ = 0.60 mm⁻¹, F(000) = 2040, crystal size 0.15 × 0.10 × 0.05 mm³, θ range for data collection: 3.43–22.97°, index ranges *hkl*: -22 to 22, -22 to 22, -13 to 13, reflections collected = 29 478, independent reflections = 3668 ($R_{\text{int}} = 0.103$), reflections with $I > 2\sigma(I) = 2766$, completeness to $\theta =$ 22.97° (99.3%), $T_{\text{max, min}} = 0.960$ and 0.833, refinement method: full-matrix least squares on F^2 , data/restraints/parameters = 3668/0/263, goodness-offit on $F^2 = 1.105$, final R indices $[I > 2\sigma(I)]$ R1 = 0.067, wR2 = 0.151, R indices (all data): R1 = 0.098, wR2 = 0.165, Largest diff. peak and hole: 0.84 and -0.65 e Å⁻³. The molecule lies on a twofold rotation axis. Data collection KappaCCD, program package WinGX, absorption correction MULTISCAN. Refinement using SHELXL-97. CCDC 623568. For crystallographic data in CIF or other electronic format see DOI: 10.1039/b614765h

¶ *Computational*: The geometries were optimized with the GAUSSIAN program using the B3LYP functional and the $6-31++G^*$ basis sets for all atoms. For the comparison of the different spin-state energies of **2** the SCF energies were used.²²

 \parallel *Fragment calculations*: The fragment calculations were performed with the Amsterdam Density Functional (ADF) program package using the BP86 functional and TZP basis sets for all atoms.²³ In the fragment calculations the molecular orbitals of the molecule is constructed as a linear combination of the orbitals of the chosen fragments. The fragments retain the geometry they have in the whole molecule and are in a restricted spin states due to prior restricted calculations.

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